

Mutli-field Open Inflation and the Effect of Interaction



京都大学 基礎物理学研究所

D1 杉村 和幸



Contents

- Motivations
- Multi-field open inflation
 - Scenario
 - Formulation and Model
 - Quantum tunneling
 - Evolution after tunneling
- Interaction effect on decay rate
- Conclusions and Discussions



Motivations



Slow-roll inflation and beyond

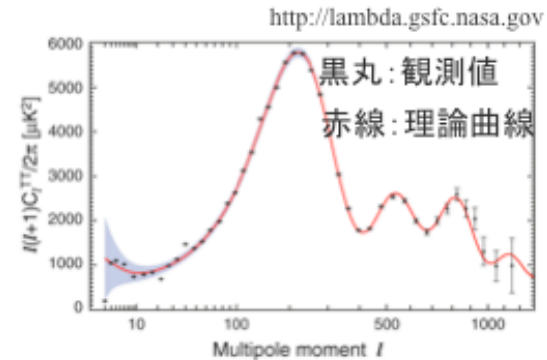
□ Slow-roll inflation is satisfactory

- Solution of
 - Horizon problem
 - Flatness problem
 - Monopole problem

- Prediction of power spectrum

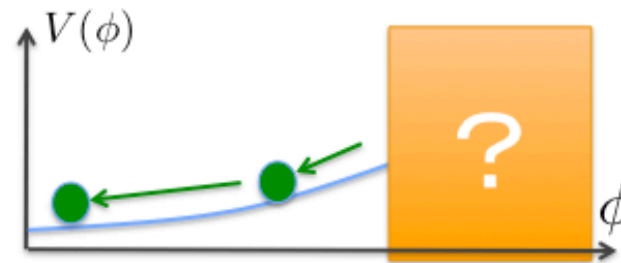


very good agreement!!



□ Slow-roll inflation is unsatisfactory

- How is it implemented in particle physics?
- How does it start?





String landscape

- String theory can implement a setting for slow-roll inflation

- scalar field

- flat potential

High dimensionality

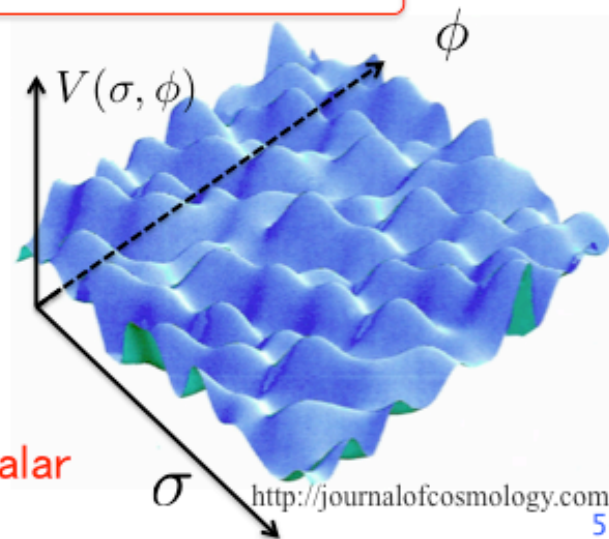
Integration out of high energy physics

many physical degrees of freedom

Non-trivial structure of vacuum

- String theory implies

- multi-field
- quantum tunneling



We will study inflation with multi-scalar field and quantum tunneling



Multi-field open inflation



Scenario

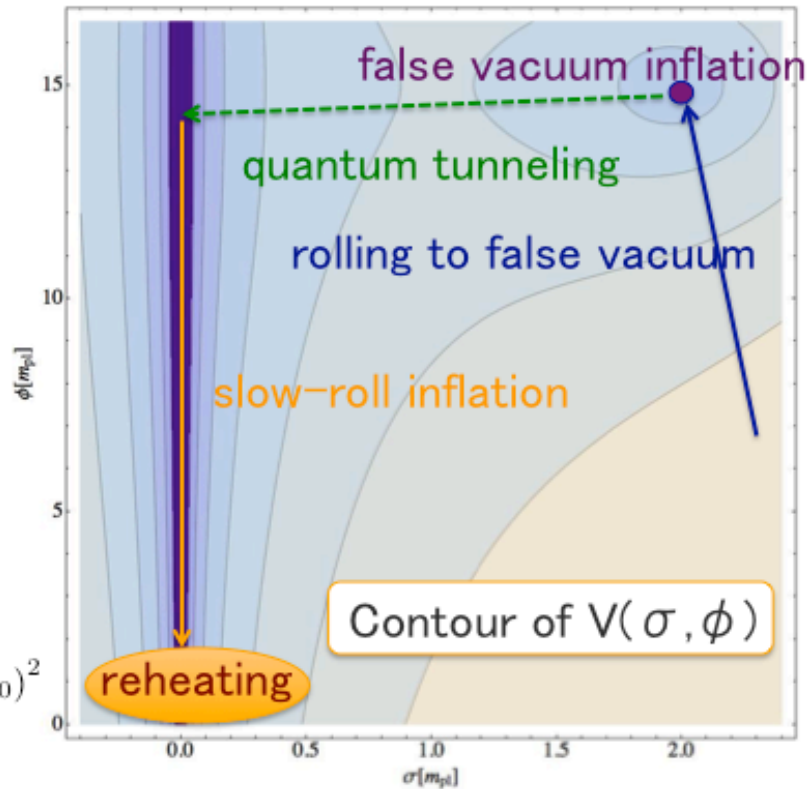
1. rolling to false vacuum
2. false vacuum inflation
3. quantum tunneling
4. slow-roll inflation
5. reheating

$$V(\sigma, \phi) = V_{\text{tun}}(\sigma) + V_{\text{infl}}(\sigma, \phi)$$

$$V_{\text{tun}}(\sigma) = \alpha \sigma^2 \left\{ (\sigma - \sigma_0)^2 + M_V^2 \right\}$$

$$V_{\text{infl}}(\sigma, \phi) = \frac{1}{2} m_\phi^2 \phi^2 + \frac{\beta}{2} \sigma^2 (\phi - \phi_0)^2$$

$$\begin{aligned} \alpha &= 0.1 & \beta &= 0.01 \\ \sigma_0 &= 2m_{pl} & m_\phi &= 10^{-6}m_{pl} \\ M_V &= 0.2m_{pl} & \phi_0 &= 15m_{pl} \end{aligned} \quad m_{pl}^2 = \frac{1}{8\pi G}$$





Review of CDL instanton

(tunneling = 1st order transition = bubble nucleation)

Single scalar field tunneling with gravity



Coleman-De Luccia instanton(1980)

under O(4)-symmetry

- Euclidean metric $ds_E^2 = dt^2 + a^2(t) (d\chi^2 + \sin^2 \chi d\Omega^2)$
- Euclidean action $S_E[\sigma, a] = 2\pi^2 \int dt a^3 \left[\frac{1}{2} \dot{\sigma}^2 + V(\sigma) - \frac{3}{m_{pl}^2} \left(\frac{\dot{a}^2}{a^2} + \frac{1}{a^2} \right) \right]$

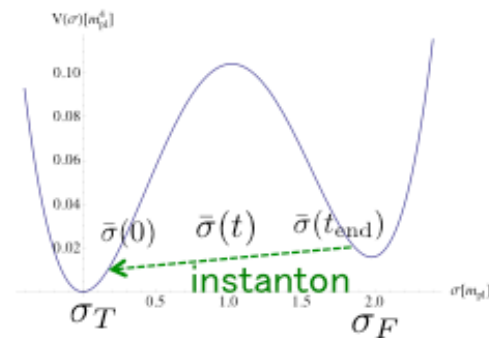


non-trivial classical path

instanton $\bar{\sigma}(t), \bar{a}(t)$

$$\bar{a}(0) = \bar{a}(t_{\text{end}}) = 0$$

$$\bar{\sigma}(0) \sim \sigma_T \quad \bar{\sigma}(t_{\text{end}}) \sim \sigma_F$$



- decay rate $\frac{(\text{tunneling probability})}{(\text{unit volume})(\text{unit time})} = O(V(\sigma)) \times \exp(-S_E[\bar{\sigma}, \bar{a}] + S_E[\sigma_F, a_F])$
 - instanton
 - trivial path at false vacuum
- evolution inside bubble
 - metric $ds^2 = -dt^2 + a^2(t) (dr^2 + \sinh^2 r d\Omega^2)$ open universe
 - initial conditions (at t=0)
 - $\sigma(0) = \bar{\sigma}(0) \quad a(0) = 0$
- evolution is same as the usual open universe



Multi-field tunneling with gravity

- Multi-field extension of Coleman-De Luccia instanton
 - instanton $\bar{\sigma}(t), \bar{\phi}(t), \bar{a}(t)$

instanton is a solution
for these equations \rightarrow

EOMs

$$\begin{aligned} \frac{\ddot{\bar{a}}}{\bar{a}} + \frac{1}{3m_{pl}^2} (\dot{\bar{\sigma}}^2 + \dot{\bar{\phi}}^2 + V(\bar{\sigma}, \bar{\phi})) &= 0 \\ \ddot{\bar{\sigma}} + 3\frac{\dot{\bar{a}}}{\bar{a}}\dot{\bar{\sigma}} - V_{\bar{\sigma}}(\bar{\sigma}, \bar{\phi}) &= 0 \\ \ddot{\bar{\phi}} + 3\frac{\dot{\bar{a}}}{\bar{a}}\dot{\bar{\phi}} - V_{\bar{\phi}}(\bar{\sigma}, \bar{\phi}) &= 0 \end{aligned}$$

boundary conditions

$$\begin{aligned} \bar{a}(0) = \bar{a}(t_{\text{end}}) &= 0. \\ \dot{\bar{\sigma}}(0) = \dot{\bar{\sigma}}(t_{\text{end}}) &= 0 \\ \dot{\bar{\phi}}(0) = \dot{\bar{\phi}}(t_{\text{end}}) &= 0 \\ \dot{\bar{a}}(0) = -\dot{\bar{a}}(t_{\text{end}}) &= 1 \end{aligned}$$

- decay rate

$$\frac{(\text{tunneling probability})}{(\text{unit volume})(\text{unit time})} = O(V(\sigma, \phi)) \times \exp(-S_E[\bar{\sigma}, \bar{\phi}, \bar{a}] + S_E[\sigma_F, \phi_F, a_F])$$

- initial conditions inside bubble

$$\sigma(0) = \bar{\sigma}(0) \quad \phi(0) = \bar{\phi}(0) \quad a(0) = 0$$



evolve as an open Friedmann universe

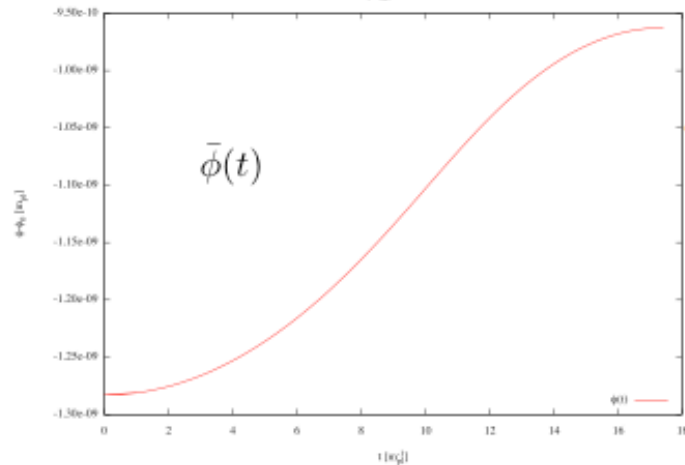
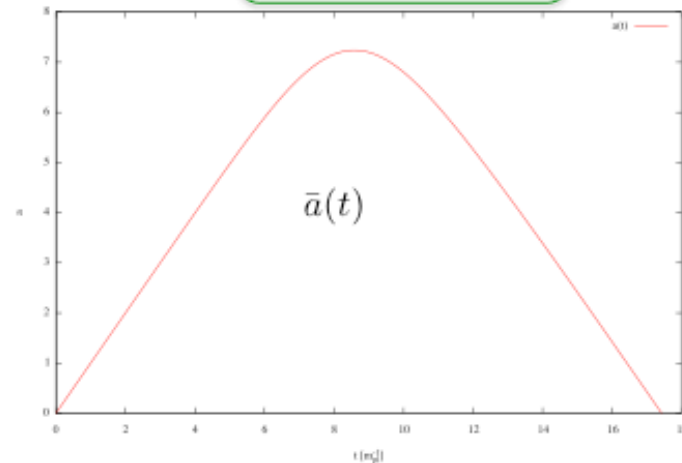
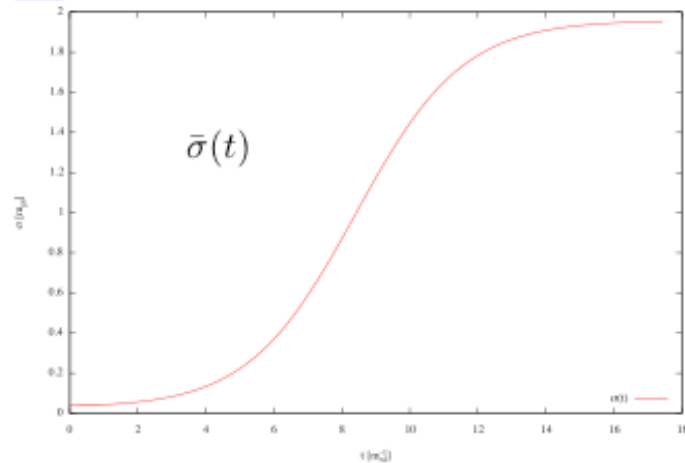
We will

1. construct a multi-field instanton
2. develop the universe inside bubble



Quantum tunneling

$$\begin{aligned} \alpha &= 0.1 & \beta &= 0.01 \\ \sigma_0 &= 2m_{pl} & m_\phi &= 10^{-6}m_{pl} \\ M_V &= 0.2m_{pl} & \phi_0 &= 15m_{pl} \end{aligned}$$



solution for EOMs and boundary conditions

• numerically obtained multi-field instanton

• inflaton ϕ moves but a little

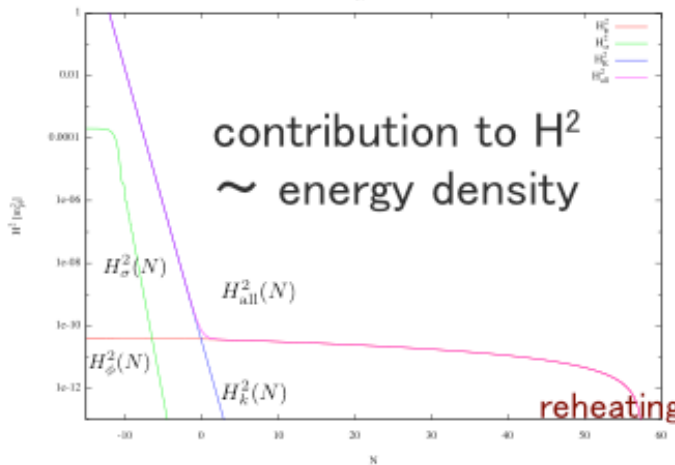
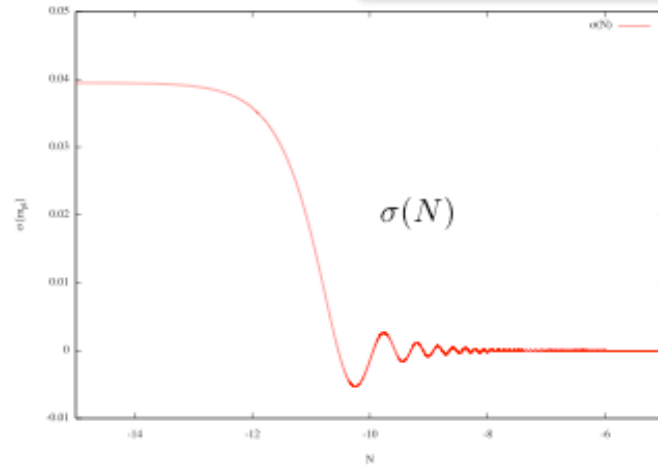
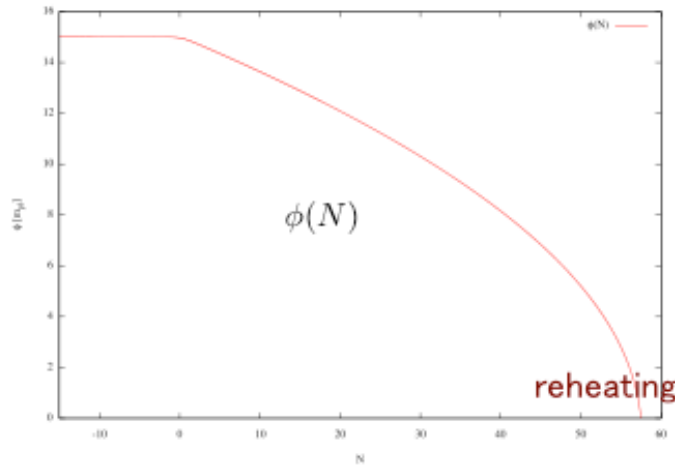
• initial condition for a nucleated bubble

$$\sigma(0) = \bar{\sigma}(0) \quad \phi(0) = \bar{\phi}(0) \quad a(0) = 0$$



Evolution after tunneling

$\alpha = 0.1$	$\beta = 0.01$
$\sigma_0 = 2m_{pl}$	$m_\phi = 10^{-6}m_{pl}$
$M_V = 0.2m_{pl}$	$\phi_0 = 15m_{pl}$



- initial conditions
 $\sigma(0) = \bar{\sigma}(0) \quad \phi(0) = \bar{\phi}(0) \quad a(0) = 0$

- e-folds
 $N = \ln(a(t)/a_i)$ scale factor at the beginning of inflation

starting from curvature dominant universe

➡ potential energy win against the curvature

➡ slow-roll inflation for ~ 60 e-folds

Realization of open inflation scenario!!



Interaction effect on decay rate



decay rate and interaction effect

Q. Does an interaction effect make decay rate larger, or smaller?

- potential

$$V(\sigma, \phi) = V^{(0)}(\sigma) + \underbrace{V_{\text{int}}(\sigma, \phi)}_{\text{interaction part}} \quad (V_{\text{int}}(\sigma, \phi_F) = 0)$$

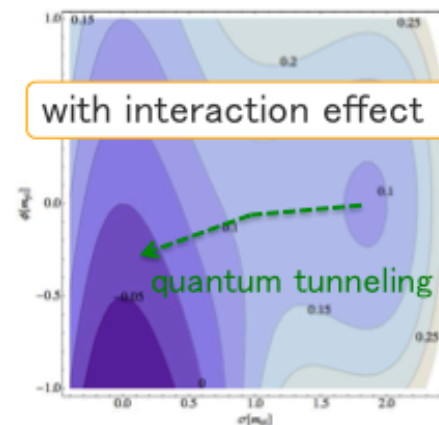
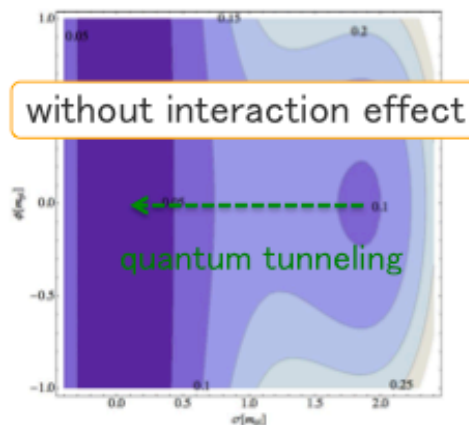
- decay rate

$$\frac{(\text{tunneling probability})}{(\text{unit volume})(\text{unit time})} = O(V(\sigma, \phi)) \times \exp(-S_E[\bar{\sigma}, \bar{\phi}, \bar{a}] + S_E[\sigma_F, \phi_F, a_F])$$

$$\frac{\partial V_{\text{int}}(\sigma, \phi_F)}{\partial \phi} = 0 \quad \text{without interaction effect}$$

$\rightarrow \bar{\phi}(t) = \phi_F$

changing V_{int} with $V^{(0)}$ fixed $\rightarrow V_{\text{int}}$ dependence of decay rate





Decay rate dependence on interaction

- potential ansatz

$$V^{(0)}(\sigma) : \text{fixed} \quad V_{\text{int}}(\sigma, \phi) = \begin{cases} \lambda(\phi - \phi_F) & \sigma \sim \sigma_T \\ \frac{m_{\phi F}^2}{2}(\phi - \phi_F)^2 & \sigma \sim \sigma_F \end{cases}$$

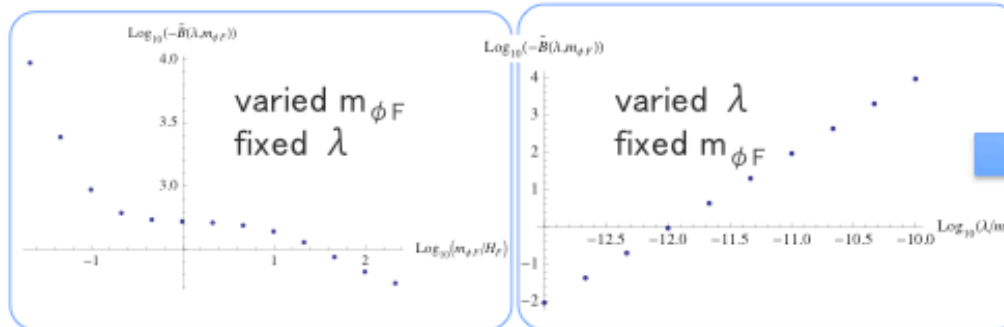
- indicator of decay rate

$$\tilde{B}(m_{\phi F}, \lambda) = -(-S_E[\bar{\sigma}, \bar{\phi}, \bar{a}] + S_E[\sigma_F, \phi_F, a_F])|_{m_{\phi F}, \lambda}$$

$$+ \frac{(-S_E[\bar{\sigma}, \bar{\phi}, \bar{a}] + S_E[\sigma_F, \phi_F, a_F])|_{m_{\phi F}, \lambda=0}}{\text{action without interaction effect (independent of } m_{\phi F})}$$

$\tilde{B} \downarrow \rightarrow$ decay rate \uparrow

- results



$$\tilde{B} < 0$$

$\tilde{B} \downarrow \leftarrow \lambda \uparrow$
 $\tilde{B} \downarrow \leftarrow m_{\phi F} \downarrow$

A. The interaction effect make decay rate larger!!

interpretation : tunneling to a lower energy state



Conclusions and discussions



Conclusions

- We studied about Multi-field open inflation, which is motivated by string landscape
- The Coleman De Luccia instanton method was extended to the multi-field case
- For a certain potential, a multi-field instanton can be obtained and open inflation is realized
- The effects of interaction make decay rate larger because of tunneling to a lower energy state



Discussions

- Initial state of the universe is different from Bunch-Davies vacuum because of tunneling
- Calculation of Non-Gaussianity from this effect and comparison with observation such as PLANCK is a future work

- Tunneling field experiences damped oscillation after tunneling
- Interaction with tunneling field may cause some features in inflaton power spectrum