

# CMB lensing-Galaxy 相互相関を用いた Primordial Non-Gaussianity への制限

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## Overview

The primordial non-Gaussianity (NG) affects the clustering of dark matter halo through the **scale-dependent bias**.

**Recent results**: Observations & Forecasts (1σ error)

**Obs.**  $f_{NL} = 53 \pm 25$  &  $f_{NL} = 47 \pm 21$ : from NVSS & SDSS DR6 QSOs data (Xia et al. 2010)

**For.**  $\Delta f_{NL} \sim 1-5$ : cluster counts for DES-like survey (Cunha et al. 2010)

**For.**  $\Delta f_{NL} \sim \text{few}$ : CMB Bispectrum with ideal CM experiment

CMB lensing is a powerful tool to explore the large scale structure, which can get matter distribution without uncertainty of bias.

Cross correlation between galaxy & CMB lensing can be break some degeneracy of NG and bias.

## 1. Scale-dependent bias (Dale et al. 2008, Slosar et al. 2008)

Deviations from Gaussian initial conditions are commonly parameterized in terms of the dimensionless  $f_{NL}$  parameter

### Primordial non-Gaussianity (NG) of the local type

$$\Phi = \phi + f_{NL}(\phi^2 - \langle \phi^2 \rangle) \quad (\text{Komatsu & Spergel 2001})$$

The effect of the primordial NG of the local type is seen in the clustering of halos through a scale-dependent bias

### Galaxy power spectrum

Gaussianity non-Gaussianity

$$P_g(k) = b_0^2 P(k) \rightarrow [b_0 + \Delta b(k)]^2 P(k)$$

$$\Delta b(k) = \frac{3(b_0 - 1)f_{NL}\Omega_m H_0^2 \delta_c}{D(z)k^2 T(k)}$$

$b_0$ : linear bias  
P(k): matter power spectrum  
D(z): growth factor  
T(k): transfer function

Primordial NG of the local type gives rise to a strong scale-dependent bias on large scales ( $\propto k^{-2}$ ), while the bias is roughly constant in the Gaussian case.

## 2. Galaxy-CMB lensing cross correlation signal

The cross correlations between the CMB and the galaxy are well known as providing additional information other than their respective autocorrelation.

We introduce the cross correlation between the CMB lensing and the galaxy angular distribution to estimate errors in constraining cosmological parameters.

### cross correlation angular power spectrum

$$C_l^{\psi g} = \frac{2}{\pi} \int k^2 dk P(k) \Delta_l^{\psi}(k) \Delta_l^g(k)$$

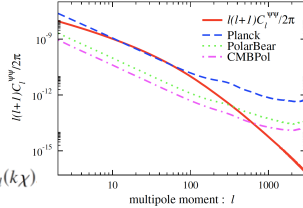
galaxy distribution (counts)

$$\Delta_l^g(k) = \int dz \frac{dN}{dz} b(k, z) T(k) D(z) j_l(k\chi)$$

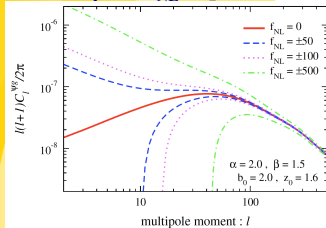
CMB lensing potential

$$\Delta_l^{\psi}(k) = -2 \int_0^{z_*} d\chi T_{\Psi}(\chi; \eta_0 - \chi) \left( \frac{\chi_* - \chi}{\chi_* X} \right) j_l(k\chi)$$

### $C_l^{\psi\psi}$ & Noise spectrum



### $C_l^{\psi g}$ & $f_{NL}$ dependence



Lensing potential  $\psi$  can be reconstructed from lensed T, E, B. (Hu & Okamoto 2002)

Planck is not sensitive to CMB lensing so much.

NG rises the power on large scale  
For the high-z and highly biased objects, the effect of NG appears more pronouncedly.

## 3. Fisher Information Matrix (Tegmark et al. 1997)

Fisher matrix

$$F_{ij} = \sum_{l=2}^{l_{\max}} \sum_{XX', YY'} \frac{\partial C_l^{XX'}}{\partial \theta_i} (\text{Cov}_l^{-1})_{XX'YY'} \frac{\partial C_l^{YY'}}{\partial \theta_j}$$

$\text{Cov}_l$ : Covariance matrix ( $XX', YY' \in TT, EE, TE, \psi\psi, T\psi, gg, Tg, \psi g$ )

To include  $\psi g$ , we calculate  $\text{Cov}_l$  as

$$\text{Ex. } \text{Cov}_l = \frac{2}{(2l+1)f_{\text{sky}}} \begin{bmatrix} (C_l^{TT})^2 + 2(C_l^{TE})^2 [C_l^{\psi\psi} (C_l^{Tg})^2 + C_l^{gg} (C_l^{T\psi})^2 - 2C_l^{T\psi} C_l^{Tg} C_l^{\psi g}] \\ 2(C_l^{\psi g})^2 - C_l^{\psi\psi} C_l^{gg} \end{bmatrix}$$

$$\text{Cov}_l = \frac{2}{(2l+1)f_{\text{sky}}} \begin{bmatrix} (C_l^{\psi\psi})^2 + C_l^{\psi\psi} C_l^{gg} \\ \psi g \psi g \end{bmatrix}$$

Marginalized 1σ error:  $\sigma(\theta_i) = \sqrt{(F^{-1})_{ii}}$

## 4. Result: Signal to Noise

Compare the S/N of some cross-correlations

$$\left(\frac{S}{N}\right)^2 = f_{\text{sky}} \sum_2^{l_{\max}} (2l+1) \frac{(C_l^{XY})^2}{(C_l^{XX})^2 + (C_l^{YY})^2 + N_l^{XX}(C_l^{YY})^2 + N_l^{YY}}$$

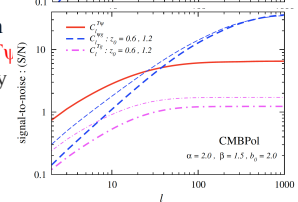
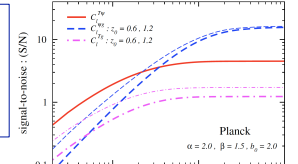
### Survey parameters

$f_{\text{sky}}$ : sky coverage

for galaxy survey ( $C_l^{\text{gg}}, C_l^{\text{Tg}}, C_l^{\text{wg}}$ ) -  $f_{\text{sky}} = 0.10$

for CMB ( $C_l^{\text{TT}}, C_l^{\text{EE}}, C_l^{\text{TE}}, C_l^{\psi\psi}, C_l^{\text{T}\psi}$ ) -  $f_{\text{sky}} = 0.65$

$N_g = 10^6$ : total number of galaxies



S/N does not increase in high- $l_{\max}$ .

=> dominated by noise in small scale region

For high- $l_{\max}$ ,  $\psi g$  get larger S/N than Tg or Tψ

$\psi g$  will get larger S/N by future CMB survey which is more sensitive to CMB lensing (e.g. CMBPol).

$\psi g$  can be an important observation value !!

## 5. Result: Parameter forecast (for Planck or CMBPol)

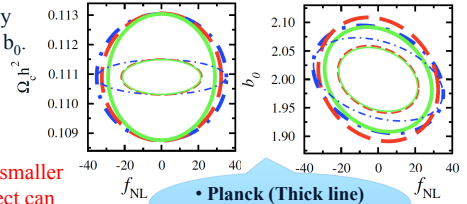
To see the contribution of  $\psi g$  for constraining NG, compare the 3 cases.

- Case I:  $C_l^{\text{TT}}, C_l^{\text{EE}}, C_l^{\text{TE}}, C_l^{\psi\psi}, C_l^{\text{T}\psi}, C_l^{\text{gg}}, C_l^{\text{Tg}}, C_l^{\text{wg}}$  (without  $C_l^{\text{Tg}}, C_l^{\text{wg}}$ )
- Case II:  $C_l^{\text{TT}}, C_l^{\text{EE}}, C_l^{\text{TE}}, C_l^{\psi\psi}, C_l^{\text{T}\psi}, C_l^{\text{gg}}, C_l^{\text{Tg}}, C_l^{\text{wg}}$  (without  $C_l^{\text{Tg}}, C_l^{\text{T}\psi}$ )
- Case III:  $C_l^{\text{TT}}, C_l^{\text{EE}}, C_l^{\text{TE}}, C_l^{\psi\psi}, C_l^{\text{T}\psi}, C_l^{\text{gg}}, C_l^{\text{Tg}}, C_l^{\text{wg}}$  (full)

$f_{NL}$  degenerates especially with linear bias parameter  $b_0$ .

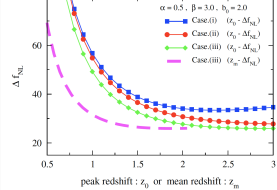
However,  $f_{NL}$  does not degenerate with other cosmological parameters so much.

The error of  $f_{NL}$  become smaller by including  $\psi g$ . This aspect can be seen more clearly for CMBPol.



Planck (Thick line)  
CMBPol (Thin line)

### redshift dependence of $\Delta f_{NL}$



The error of  $f_{NL}$  depends extensively on the peak redshift of the galaxy distribution  $z_0$ .

=> the effect of NG become large with  $z_0$ .  
For large  $z_0$ , the range of  $\Delta f_{NL}$  is small.  
=> the amplitude of the matter density is small.

## for HSC-like survey

HSC: Hyper-Suprime Cam (Subaru Telescope)

Assuming finite observation time, survey parameters are related as below.

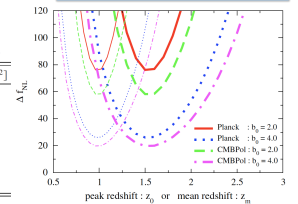
### Survey parameters

$$z_m = 0.9 \left( \frac{t_{\text{exp}}}{30 \text{ min}} \right)^{0.067}$$

$$n_L = 35 \left( \frac{t_{\text{exp}}}{30 \text{ min}} \right)^{0.44} \text{ [deg}^{-2}\text{]}$$

$$\text{area} = \pi \left( \frac{\text{field of view}}{2} \right)^2 \frac{T_{\text{total}}}{1.1 \times t_{\text{exp}} + t_{\text{op}}}$$

$z_0$	$f_{\text{sky}}$	$n_L$ [deg <sup>-2</sup> ]
1.0	0.41	0.8
1.3	0.39	4.7
1.4	0.34	7.6
1.5	0.26	12.0
1.6	0.16	18.3
1.7	0.09	27.2
1.9	0.02	56.5
2.2	0.002	148.1



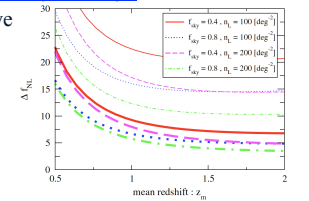
## for future wide field galaxy surveys

Future galaxy surveys will be able to observe the wider field and more objects.

=> get larger  $f_{\text{sky}}$  &  $N_g$  (or  $n_L$ ).

Some combinations of the survey parameters ( $f_{\text{sky}}, n_L$  &  $z_0$ ) achieve  $\Delta f_{NL} < 5$ .

This results are competitive to the CMB Bispectrum or Trispectrum constraints.



## 6. Summary & Conclusion

Cross correlation  $\psi g$  plays an important role to break some degeneracy between  $f_{NL}$  and  $b_0$ , and to constrain on  $f_{NL}$  more tightly.

It is competitive constraint on  $f_{NL}$  CMB higher-order correlation (Bispectrum, Trispectrum, etc...) with ideal CMB observation.

The constraint through the scale-dependent bias is sensitive only to the local-type NG. For the constraints on the other NG models, we must consider higher-order correlation.